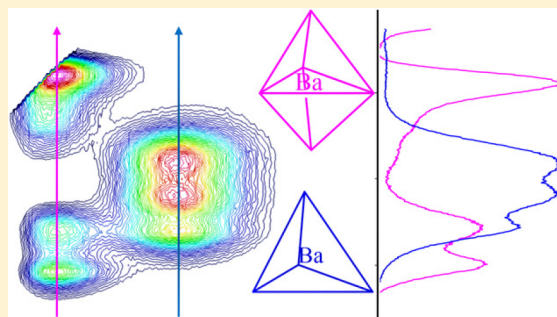


Luminescence of Atomic Barium in Rare Gas Matrices: A Two-Dimensional Excitation/Emission Spectroscopy Study

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ABSTRACT: A detailed characterization is made of the distinct sites occupied by atomic barium isolated in the three rare gas hosts Ar, Kr, and Xe in excitation scans extracted from the recorded total $6s6p\ ^1P_1 \rightarrow (6s)^2\ ^1S_0$ fluorescence. Extensive use has been made of two-dimensional excitation/emission (2D-EE) spectroscopy to achieve a comprehensive characterization for the wide variety of sites present in the Ba/RG matrix systems. The 2D-EE technique has proved to be a very powerful method to probe the effects of strong intersite reabsorption when extensive spectral overlap occurs between emission and resonance $6s6p\ ^1P_1 \leftarrow (6s)^2\ ^1S_0$ absorption of barium atoms occupying multiple sites. Two-dimensional excitation/emission scans have also been used in this study to monitor the effects of sample annealing and thereby identify the thermally stable sites of isolation. Sites of the same type occupied by atomic barium in the three host solids are identified in resolved excitation spectra and are associated on the basis of the observed matrix shift versus host polarizability. Following site associations, the photophysical properties of each matrix site were characterized revealing that the Stokes shift was greatest in the blue site, smallest for the violet site, and intermediate for the green site. The emission temperature dependences and excited state lifetimes were recorded, indicating that measured radiative lifetimes of 4–5 ns were in good agreement with the gas phase value of 8.4 ns when corrected for the effective field of the solids. The only exception to this was the blue site in Ba/Xe, where a nonradiative quenching channel exists even at 9.8 K that competes effectively with the nanosecond fluorescence. An unusual, asymmetric $2 + 1$ excitation band has been recorded for atomic barium in the three rare gas hosts in addition to the threefold split, Jahn–Teller bands typically observed for $P \leftarrow S$ absorptions of matrix-isolated metal atoms. Possible assignments of the sites responsible for these band shapes are made on the basis of recent spectral simulations obtained from molecular dynamics calculations on the Ba/Xe system.



I. INTRODUCTION

This article presents an overview of the total luminescence recorded for barium atoms isolated in solid Ar, Kr, and Xe and a detailed examination of the distinct excitation bands arising from the variety of sites occupied by this guest metal atom in the three solid rare gases. With the inclusion of data for the Ba/Ar and Ba/Kr systems, the present contribution expands the matrix emission data and, in so doing, provides a more comprehensive analysis of the site attributions published by us recently for Ba/Xe.¹ The first Ba/rare gas (RG) matrix luminescence spectra were reported more than three decades ago by Balling and Wright² (B&W), while results for Ba/Ar and Ba/Xe were published recently by Fairbank and co-workers.³ B&W's emission spectra are, however, not consistent with those recorded in the present study as shown by the red traces in Figure 1, even though both were produced with photoexcitation of the dominant $(6s6p)\ ^1P_1$ absorption bands. The most significant difference is that our Ba/Ar spectrum shows a pair of resolved emission bands located at 537 and 550 nm while that published in 1985 by Balling and Wright² consisted of a single, longer wavelength band at approximately 568 nm. The reason for the pronounced difference is, as will be shown in this work, the matrix samples prepared in B&W's study were strongly absorbing which had the effect that complete

reabsorption of the blue emission features occurred leaving only emission from the band furthest to the red. The authors B&W acknowledged that the observed emission probably originated from a single site, with the intensity of this band enhanced by the successive reabsorptions of higher energy emission bands. In an effort to identify all of the site-related Ba/RG emission bands—especially the emission features furthest to the blue—samples having low optical densities ($OD < 0.5$) were prepared and the total 1P_1 state luminescence was recorded with two-dimensional excitation–emission (2D-EE) spectroscopy. Fortunately, the deposition conditions required to achieve a high degree of atomic isolation result in relatively weakly absorbing samples. In addition, strong matrix annealing was used to further simplify the emission spectra by removing thermally unstable sites, thereby reducing the number of features present in the complex absorption spectra.

Based on absorption data, it was proposed by us in an earlier paper⁴ that three thermally stable sites of isolation exist for Ba in solid Ar and Kr and at least two sites occur in Xe. This

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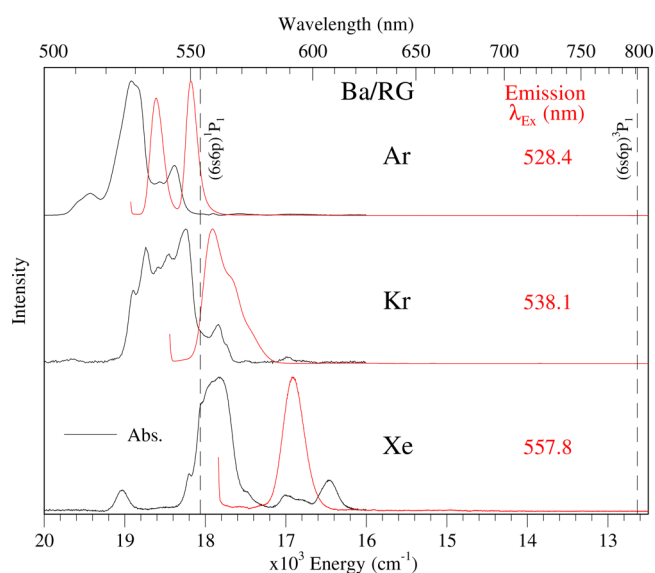


Figure 1. Absorption and emission spectra recorded at 10 K for annealed Ba/RG samples. The black traces are the absorption scans recorded after sample annealing to temperatures of 28, 38, and 60 K for Ar, Kr, and Xe, respectively. Exciting into the absorption band maxima, at the indicated wavelengths, produces intense emission in each RG host as shown by the red traces. The position¹⁰ of the atomic barium $(6s6p) \ ^1P_1 \leftarrow (6s^2) \ ^1S_0$ transition at 553.7 nm in the gas phase, is shown by the dashed vertical line, as is the $(6s6p) \ ^3P_1 \leftarrow (6s^2) \ ^1S_0$ transition at 791.35 nm. The structures evident in the absorption spectrum of Ba/Xe on both sides of the atomic $^1P_1 \leftarrow ^1S_0$ band are due to barium dimer. Their concentration dependence and emission characteristics have been analyzed and presented in ref 4 for all three host solids.

analysis assumed the occurrence of a Jahn–Teller threefold split band for each matrix site, as was the case for the other spherically symmetric M/RG systems studied by the Maynooth Group.^{5,6} However, the present analysis of the $(6s6p) \ ^1P_1 \leftrightarrow (6s^2) \ ^1S_0$ luminescence of Ba isolated in the three RG solids (Ar, Kr, and Xe) reveals a more complex situation. Excitation spectroscopy was particularly useful in revealing bands severely overlapped in the absorption spectra, the details of which will now be presented.

The approach adopted in the present study to interpret the large variety of site-related excitation (absorption) bands relies on an observation first presented by Laursen and Cartland,⁷ who showed that the matrix shifts for atomic $P \leftarrow S$ absorptions of the group 12 metal atoms are approximately linear when plotted against the polarizability of the RG hosts. The basis for this simple dependence only became evident when the molecular constants of the 1:1 van der Waals complexes formed between the group 12 metal atoms Zn, Cd, and Hg and the heavier rare gas solids Ar, Kr, and Xe were obtained. All these complexes are now known⁸ to have relatively short M–RG bond lengths that closely match the nearest neighbor distances in the rare gas solids. The good size match of the guest atoms and the host vacancies favors single substitutional site occupancy for all three metal atoms in the heavier rare gas solids, Ar–Xe. With the same (single vacancy) site occupied in these M/RG systems, a linear dependence of matrix shift with host polarizability is observed. This behavior was also found for atomic Na,⁹ Mn,⁵ and Eu⁶ even though considerable guest/host size mismatches exist, resulting in more complex absorption profiles from the occupancy of these

guest atoms in multiple sites. In these cases, polarizability plots allowed for associations to be made for the same site types occupied by metal atoms in the three rare gas solids. Accordingly, matrix bands that exhibit a linear dependence of the matrix shift versus rare gas polarizability are expected for metal atoms trapped in a particular site type. On this basis, the same site occupied by atomic barium in the three solid rare gases can be identified from plots of the matrix shifts versus polarizabilities of the host solids.

State assigning the emission bands shown in Figure 1 is simple in that, based on their proximity to the gas phase position¹⁰ (indicated by the dashed vertical line), the observed emissions are naturally attributed to the $6s6p \ ^1P_1 \rightarrow 6s^2 \ ^1S_0$ transition of atomic Ba. Scans made out to 950 nm (the limit of our detection system) did not reveal any relaxed emission from the lower energy 1D_2 or 3P_1 states of atomic barium. Although our experimental setup does not allow detection of the 3D_1 state (which for the free atom occurs in the NIR at 1170 nm), the absence of the nearby metastable states and the low intensity laser excitation used indicate negligible involvement of transitions from metastable levels. This observation is consistent with similar findings by Lebedev et al.¹¹ for atomic barium isolated in solid, pressurized helium when resonant excitation at approximately $18\,500 \text{ cm}^{-1}$ was used instead of high power radiation of the frequency-tripled Nd:YAG laser at $28\,195 \text{ cm}^{-1}$. Under resonant excitation of Ba/He with a weaker OPO laser, the strongest fluorescence observed was the line at $18\,115 \text{ cm}^{-1}$ which corresponds to the transition from the $6s6p \ ^1P_1$ state to the ground $6s^2 \ ^1S_0$ state. This is what we also observe with resonant excitation of atomic barium in the solid rare gases (SRGs) Ar, Kr, and Xe and accordingly assign to the $6s6p \ ^1P_1 \rightarrow 6s^2 \ ^1S_0$ transition. This assignment is confirmed with excited state lifetime measurements conducted in the present study where nanosecond decay times were recorded in each solid.

In contrast to the relatively straightforward state assignment, the emission band profiles shown in Figure 1 suggest that the luminescence of Ba in these hosts is complex. In particular, the presence of multiple emission bands in the Ba/Ar and Ba/Kr systems points to a wider variety of site occupancy than in Ba/Xe. This situation is manifested also in the absorption bands shown in Figure 1. The detailed examination presented in this study relies on the excitation spectra recorded for the numerous distinct emission bands that exist in each host, to provide resolved bands which are strongly overlapping in absorption and whose origins are the multiple site occupancy of the barium atom. In a recent combined experimental and theoretical study of the Ba/Xe system,¹ the resolved excitation spectra of the thermally stable blue and violet bands were attributed to a four-atom tetravacancy (TV) and a five-atom vacancy, grain boundary site, respectively. The minor green band present in Ba/Xe was attributed to a six-atom hexavacancy (HV) site. From a site occupancy perspective, Ba/Xe is a simpler system than either of the lighter hosts and forms a basis for the discussion of the sites occupied by barium atoms in Ar and Kr. The strength of the present study, in which the combined results of the three host solids are analyzed together, is that its overview provides greater scope to check the site proposals made recently in the Ba/Xe system.¹

The results are presented in the following manner. First, an illustration of the effects of reabsorption, as monitored with 2D-EE spectroscopy, is provided using the Ba/Ar system as a case study. Second, an overview of the entire $(6s6p) \ ^1P_1$ state

luminescence of each Ba/RG (RG = Ar, Kr, and Xe) system is provided with the 2D-EE spectra. Only a summary of the Ba/RG matrix spectra is presented here to allow comparisons to be made between the sites present in the three systems. Full spectroscopic details of the data recorded for the individual Ba/Ar, Ba/Kr, and Ba/Xe matrices are provided elsewhere.¹² Following this, steady state and time-resolved emission scans are presented for each of the distinct trapping sites identified. Emission scans and decay profiles are then recorded at elevated temperatures to investigate the existence of any nonradiative decay processes competing with the resonance fluorescence of Ba at a given site of isolation. Finally, a discussion of the possible sites occupied in the three Ba/RG systems is provided based on the polarizability plots and the recent theoretical work done on the Ba/Xe system.¹

II. METHODS

All the data presented here were obtained, as outlined in the preceding papers,^{1,4} from samples deposited at 10 K under conditions of low metal loading which were subsequently annealed to minimum temperatures of 28, 38, and 60 K for Ar, Kr, and Xe, respectively. The experimental arrangement used to record luminescence is the same as that described previously for the Ba/Xe¹ system and expanded on below. The only significant difference is that extensive use was made of two-dimensional excitation and emission (2D-EE) spectroscopy allowing emission bands arising from multiple site occupancy to be identified. These plots also allow for changes in the luminescence to be easily tracked, such as those arising from reabsorption effects and those that occur with matrix annealing. More significantly, the excitation traces extracted from carefully selected emission regions of the 2D-EE spectra are then used to deconvolute the absorption profiles into site-specific features.

Emission spectra were recorded perpendicular to the excitation beam with an Acton Research 0.5 m SP500i monochromator using two distinct optical arrangements.^{1,12} For steady state work the dispersed output of a tungsten lamp and a Hamamatsu R928P PMT detector, maintained at -20 °C in a Products for Research (Photocool, S600) cooled housing, were used. For nanosecond time-resolved work an intensified, time-gated charge coupled device (iCCD) detector (Andor Technologies, model iStar DH720) held at -15 °C by an integrated Peltier cooling system was used with pulsed excitation from a Q-switched Nd:YAG laser (Quantel, YG-980E-10). The third harmonic of the YAG (355 nm), operating at a repetition rate of 10 Hz, was used to pump a dye laser (Quantel, TDL-90). Coumarin 540A dye was used to produce tunable laser radiation in the 523–586 nm range for pumping the resonance $6s6p\ ^1P_1 \leftarrow 6s^2\ ^1S_0$ transition of atomic barium in the solid rare gases. Laser intensities were kept low, in the range $2\text{--}6\ \mu\text{J}/\text{mm}^2$ (recorded at the sample window with a Moletron Power-max 500 A meter and PM10V1 head), to ensure one-photon transitions of the barium atoms. The excitation fluences were obtained using just the oscillator and/or preamplifier of the dye laser and values varied in the range quoted, depending on the wavelength selected. The resulting emission was detected with the time-gated iCCD camera setting the observation interval to the minimum value of 2 ns and moving in 1 ns steps to the limit of where the emission intensity was detectable. This was less than 100 ns for all of the Ba/RG systems. From the recorded time-resolved spectra, decay curves were extracted at selected emission wavelengths.

Excited state lifetimes were obtained by fitting exponential functions (after reconvolution with the laser pulse) to the recorded emission decay curves.

The iCCD detector was also used to record time-integrated (steady state) spectra by switching to its continuous-wave mode of operation. In this arrangement, the Andor iStar software and card (CCI-010), mounted on a PC, controlled the SP-500i monochromator to select the emission region of interest. For this setup the dispersed output of the tungsten lamp was used as the light source and by scanning the excitation monochromator—an ARC SP-300i—the total luminescences of the Ba/RG matrices were recorded. These scans are presented in the form of 2D-EE spectra. In all the 2D-EE spectra shown, scattered light from the excitation source has been removed for the purposes of presentation. Removal of the excitation light produces a well-defined diagonal border on the short wavelength side of the recorded emission as shown in Figures 2–4. The remainder of the scans presented in this work are from carefully annealed samples—hence the features shown are considered to arise from thermally stable sites.

III. RESULTS AND DISCUSSION

III.a. Reabsorption Effects. During the course of this work, it was noted that the OD of Ba/RG samples greatly affected the appearance of the recorded luminescence. This effect is an optical artifact arising from strong reabsorption of the emission that occurs when there is extensive overlap with absorption bands of a fully allowed atomic resonance transition. The 2D-EE spectra presented in Figure 2, recorded in the region of the $(6s6p)\ ^1P_1$ luminescence of Ba isolated in solid Ar, illustrate this effect. The data shown in the panels were obtained from two samples both deposited at 10 K, but with differing absorption strengths. The top spectrum corresponds to a sample with an optical density (OD) of 0.51 at 528 nm, while the bottom spectrum was obtained from a less absorbing sample (OD = 0.28 at 528 nm). Three emission regions are observed in both spectra, centered at approximately 537, 551, and 562 nm; however, the relative intensities of the three bands differ greatly depending on the OD of the sample. Thus, in the stronger absorbing sample the longest wavelength emission feature dominates the luminescence in agreement with the original Ba/Ar spectra recorded by B&W, where only the 568 nm band was observed. In contrast, the emission band at 551 nm is the most intense in the less absorbing sample, consistent with the Ba/Ar spectrum shown in Figure 1.

The effects of reabsorption can be reduced further by sample annealing as illustrated in Figures 3 and 4, which present the changes recorded in the Ba/Ar and Ba/Kr systems, respectively. It is clearly evident in the two figures that the red emission features are removed in both systems. In general, the shortest wavelength emission bands are more prominent and resolved in the weaker absorbing samples which have been annealed. This is a demonstration of the requirement for low optical densities to reduce reabsorption effects and thereby obtain the complete emission spectroscopy. The full extent of the spectral overlap between absorption and emission was illustrated in the recent Ba/Ar work published by Fairbank and co-workers.³ Their Ba/Ar spectrum is similar to that shown in Figure 1, except the complete band profile of the 537 nm emission was not recorded in the earlier work.

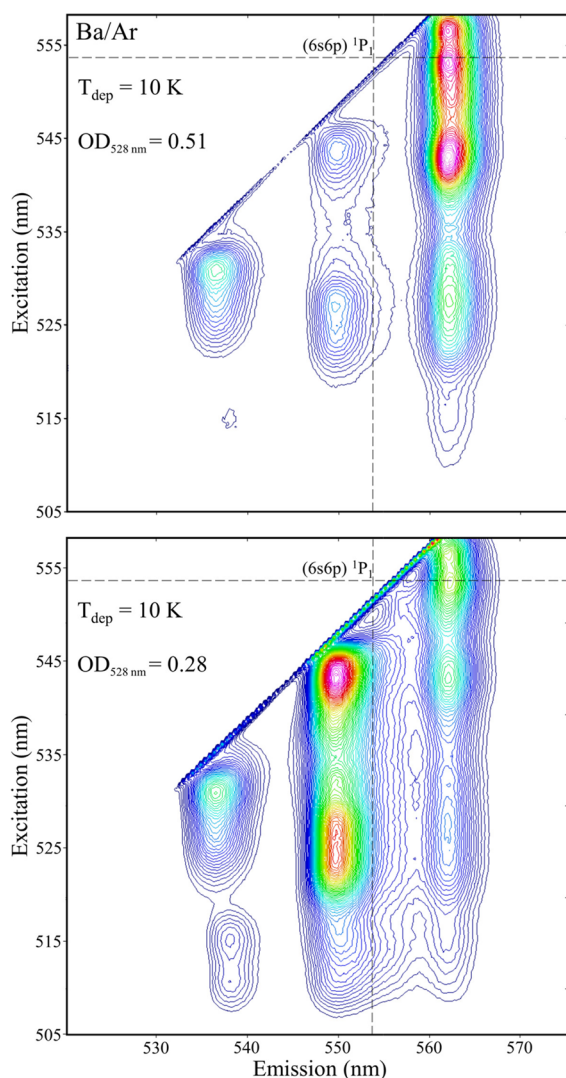


Figure 2. Two-dimensional excitation/emission (2D-EE) spectra recorded for the entire $(6s6p) \ ^1P_1$ state luminescence of Ba isolated in Ar illustrating the effects of optical density on the recorded spectra. The top panel depicts the luminescence originating from a sample with an absorption optical density (OD) of 0.51 at 528 nm. The bottom spectrum corresponds to a less absorbing sample with an OD of 0.28 (at 528 nm). Clearly the sample with the larger the OD_{528} produces much stronger red emission features with greatly diminished blue features compared with the OD = 0.28 sample. Both samples were deposited at 10 K. The gas phase position¹⁰ of the $(6s6p) \ ^1P_1 \leftarrow (6s^2) \ ^1S_0$ transition of Ba at 553.7 nm is shown, in both excitation and emission, as dashed lines.

The origins of the effects of strong reabsorption are the multisite nature of the isolation of barium atoms which produce broad and complex $(6s6p) \ ^1P_1$ absorption profiles. These absorption profiles, presented in our earlier Ba/RG work,⁴ were observed extending over the spectral range from 505 to 562 nm. Accordingly, the Ba/Ar emission bands centered at 537 and 551 nm which occur within this region are therefore very prone to competitive reabsorption. The extent of reabsorption, as shown in Figure 2, depends on the optical density of the doped RG solid. As a consequence, if a Ba/Ar sample is strongly absorbing, the lowest energy emission band dominates and the corresponding excitation spectrum will consist of a continuum spanning the entire absorption range. As each emission band corresponds to a unique site of isolation

for Ba in solid Ar, resolving site occupancy of the guest metal atom, under such conditions, will not be possible in the excitation spectra recorded. The situation described was also observed for Ba/Kr following sample deposition at 10 K. As was shown previously,¹ this effect is not as significant for Ba/Xe as the dominant emission features are Stokes shifted away from the absorption profile and Xe contains fewer sites of isolation for atomic Ba than the two lighter hosts.

III.b. Site Association. Summaries of the complete excitation spectroscopy recorded for atomic barium isolated in Ar, Kr, and Xe matrices are presented in Figure 5 and contrasted with the absorption profiles shown by the broken black traces. The excitation scans were selected by identifying the unique emission bands ($X.n$) in the 2D-EE spectra shown in Figures 3 and 4 for Ar and Kr respectively and in Figure 1 of ref 1 for Ba/Xe. It is evident that these scans provide resolved bands which are severely overlapped in the absorption spectra. Indeed many of these bands are not evident in absorption, especially in the Ba/Ar and Ba/Kr systems. It is clear in Figure 5 that the Ba/Xe system presents by far the simplest Ba/RG matrix, exhibiting just two dominant sites of isolation. In addition, a minor “green” trapping site was also identified for atomic barium in xenon based on the distinct emission it produces. The comparison provided in Figure 5 indicates that the excitation spectra recorded for Ba in Ar and Kr present some features similar to those observed in Xe. Thus, in both Ar and Kr matrices, excitation bands associated with violet, blue, and green sites of isolation were observed. A fourth matrix “red” excitation profile was also observed to be stable in Ar and unstable in Kr. Although complex, the dominant components in the red excitation traces, in line shape and emission analyses to be shown below, were attributed to two distinct trapping sites (R1 and R2) for atomic Ba. For comparative purposes, the unstable red profile of Ba/Kr is also included in the analysis.

The traces shown in Figure 5 for the three Ba/RG systems are color coded to match the site labeling scheme, whose origins are in the polarizability model—the details of which will now be presented. Gas phase to matrix frequency shifts are calculated for each of the three components of the blue, violet, and R1 (asymmetric) site bands. As the green site excitation spectra did not present any resolved structure, the band centers were used in lieu. However, as will be shown in Figure 8, the locations of the underlying three component bands were extracted in Gaussian line shape fits of the of the green site excitation profiles. Plots of these matrix shifts versus the host polarizability are presented in Figure 6.

From the approximately linear behavior exhibited, the polarizability analysis demonstrates a correspondence between the sites of a given “color” in all three hosts. For example, the dominant site in Xe, which shows a Jahn–Teller (JT) splitting (the blue trace in Figure 5), is associated with the JT band observed in Kr and the minor, poorly resolved threefold split band in Ar. Thus, it is proposed that each site “color” represents a particular type of trapping site for atomic barium in the rare gas solids. The details of the absorption/excitation bands unique to each site of isolation are presented in Table 1 for all three RG hosts. Interestingly, it is observed that the excitation components of a given site are not exactly linearly dependent on the host. For the violet, blue, and green sites, the slope of the line from Ar to Kr is greater than that from Kr to Xe. This highlights the simplicity of the polarizability model, which does not take into account the very different interactions

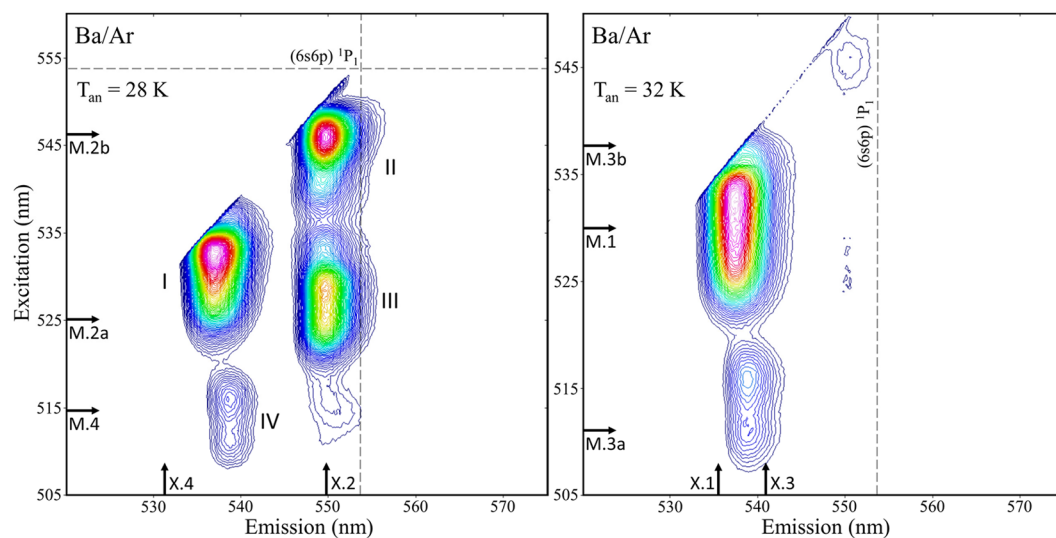


Figure 3. Two-dimensional excitation/emission spectra recorded for $(6s6p) \ ^1P_1$ state luminescence of Ba/Ar, following sample annealing. As indicated in the left panel, the emission band at approximately 562 nm has been completely removed at 28 K indicating its thermal instability, while the emission band at approximately 550 nm has almost been completely removed at 32 K. By keeping the optical density low and with sample annealing, the blue emission features can be identified and their excitation spectra recorded.

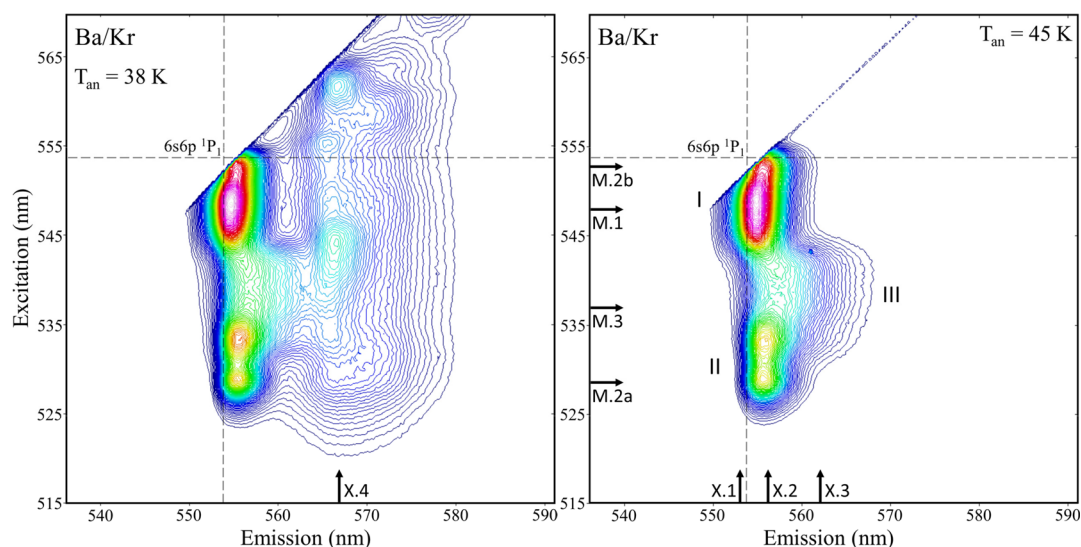


Figure 4. Effect of sample annealing on Ba/Kr system as monitored by 2D-EE spectra recorded for the $(6s6p) \ ^1P_1$ luminescence of atomic Ba isolated in krypton. As observed in the Ba/Ar system, the red emission features are completely removed with annealing to 45 K.

that occur between atomic barium and the surrounding RG environment in the ground, 1S_0 , and excited, 1P_1 , states.

III.c. Blue Site Luminescence. A summary of the blue site luminescence recorded in the three Ba/RG matrix systems is presented in Figure 7. The excitation traces of Kr and Xe both exhibit pronounced threefold splitting which is attributed to the Jahn–Teller (JT) effect, while the corresponding band in Ar seems to only show a partly resolved doublet. However, a line shape analysis conducted on this profile (shown by the three gray traces in Figure 7) reveals that a weak, higher energy band, located at approximately 513 nm, makes up the third component of the characteristic JT pattern. Based on the band splitting exhibited, the blue site of isolation must be one of high symmetry (O_h or T_d), which gives rise to a dynamic Jahn–Teller distortion in the excited P state of atomic barium.

The abundance of the blue site varies considerably depending on the host. It is the dominant site of isolation

for Ba in a Xe matrix, but in contrast, blue site occupation is reduced in Kr where the green and violet sites are similarly dominant. In Ar, this site is so weakly populated that it is difficult to detect as revealed by the poor signal-to-noise ratio evident in the excitation scan shown in the top panel of Figure 7. It seems to only form in Ar samples deposited at higher temperatures (with a very low metal flux).

Inspection of Figure 7 reveals that the expected progressive red shift of the excitation/emission (ex/em) components of this site clearly occurs in Kr to Xe as one moves to a more polarizable RG host. The behavior in Ar is not immediately evident as the blue site emission is weak (indicated by the asterisk) due to the small amount of it present in this matrix and the threefold splitting pattern in excitation is not pronounced. However, with careful analysis, the ex/em components in Ba/Ar have been identified. As indicated in Figure 7, the ex/em bands of the Ba/Ar system are located at

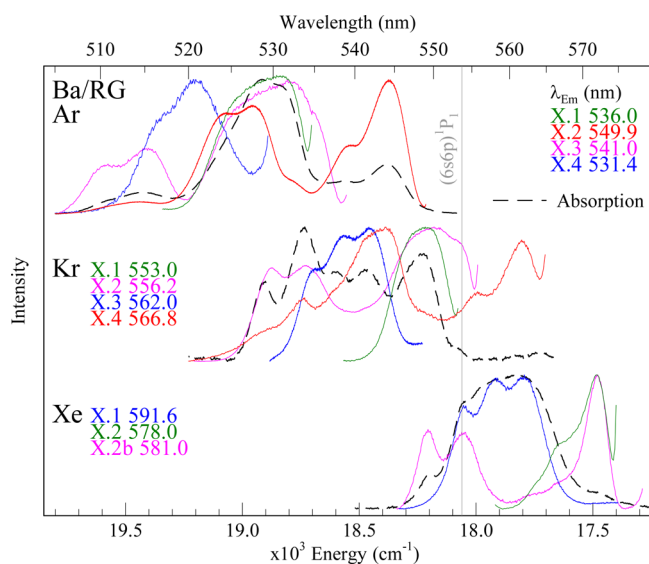


Figure 5. Site-specific $(6s6p) \ ^1P_1 \leftarrow (6s^2) \ ^1S_0$ excitation spectra extracted at the indicated emission wavelengths (X.n) in the 2D-EE scans recorded for Ba atoms isolated in Ar, Kr, and Xe matrices at 10 K. All the spectra shown were recorded following sample annealing. The absorption spectra of Ba (dashed black traces) are complicated by the presence of numerous distinct matrix sites which overlap considerably. The color of each excitation trace represents the site of isolation. Solid colored lines are used to present the excitation scans for the violet, blue, green, and red sites. For clarity, some of the excitation profiles have been truncated in an attempt to remove any overlapping components.

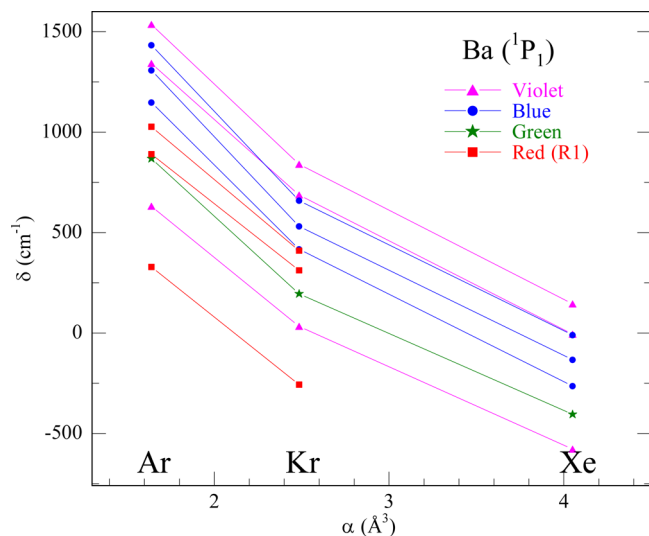


Figure 6. Plot of Ba/RG matrix frequency shifts (δ , cm^{-1}) relative to gas phase position of the $(6s6p) \ ^1P_1 \leftarrow (6s^2) \ ^1S_0$ transition of atomic Ba, observed for violet, blue, green, and red sites versus RG host polarizabilities (α , \AA^3).

the highest energy for the three rare gas hosts. Values of the matrix shifts are given in Table 2. The magnitude of the shift is related to the strength of the short-range M–RG repulsive interaction in the excited P state and strongly depends on the size of the vacancy. For instance, the most polarizable host, Xe, has the largest lattice parameter and therefore has the largest blue site radius. This leads to less confinement of the excited 6p orbital of Ba. The effects of increased host polarizability and a larger lattice vacancy both contribute to produce a moderate

Table 1. Photophysical Characteristics of Ba/RG Sites of Isolation As Revealed by Excitation Spectroscopy Recorded for the $(6s6p) \ ^1P_1 \leftarrow (6s^2) \ ^1S_0$ Transition^a

Ba/RG site	component	λ_{ex} (nm)	ν (cm^{-1})	δ (cm^{-1})
argon				
violet	1	510.3	19 596	1536
	2	515.4	19 402	1342
	3	535.0	18 692	631
blue	1	513.0	19 493	1433
	2	516.3	19 369	1308
	3	520.6	19 209	1148
green	1	528.3	18 929	868
red	1	523.9	19 088	1027
	2	527.7	18 950	890
krypton				
violet	1	529.1	18 900	840
	2	533.4	18 748	687
	3	552.7	18 093	33
blue	1	534.2	18 720	659
	2	537.9	18 591	531
	3	541.2	18 477	417
green	1	547.8	18 255	195
red ^b	1	541.4	18 471	410
	2	544.3	18 372	312
	3	561.7	17 803	−257
xenon				
violet	1	549.3	18 205	145
	2	553.9	18 054	−6
	3	572.0	17 483	−578
blue	1	554.0	18 051	−10
	2	557.8	17 928	−133
	3	561.9	17 797	−264
green	1	566.4	17 655	−405

^a δ represents the gas phase to RG matrix shift (given in units of wavenumber) calculated from the gas phase value¹⁰ for this transition at $18\,060.261 \text{ cm}^{-1}$ (553.7 nm). The wavelength uncertainties in the excitation values are estimated to be $<0.1 \text{ nm}$. In the working spectral range this translates into an uncertainty of $\pm 4 \text{ cm}^{-1}$. ^bThe red site in Kr was shown to be a thermally unstable site of isolation but is included to allow comparison with thermally stable red site in Ar.

gas-to-matrix red shift of -133 cm^{-1} for the blue site in Xe. Contrasting this, solid Ar is the least polarizable of the three hosts and also has the smallest lattice parameter. Therefore, the blue site radius in Ar will be smaller than in Xe and the excited 6p orbital of Ba will be expected to experience a greater repulsive interaction. The combination of both effects yields a large blue shift of $+1308 \text{ cm}^{-1}$ (measured from the band center) for the blue site in Ar.

The photophysical characteristics of the site-specific fluorescence bands in the three Ba/RG systems are collected in Table 2. Clearly the Stokes shift (SS) of the blue site emission, computed from the central component of the JT excitation band in each host, increases from Ar to Xe. Its value in the heaviest host (Xe, 1024 cm^{-1}) is almost twice that measured for the lightest (Ar, 543 cm^{-1}). This difference is graphically illustrated in Figure 7 by the length of the “SS lines” shown there. In addition, the emission bandwidths (full width at half-maximum, fwhm) are also observed to increase from Ar to Xe. Both of these effects indicate that the minimum of the excited $\text{Ba}(^1P_1)/\text{Xe}$ potential energy surface is displaced, with

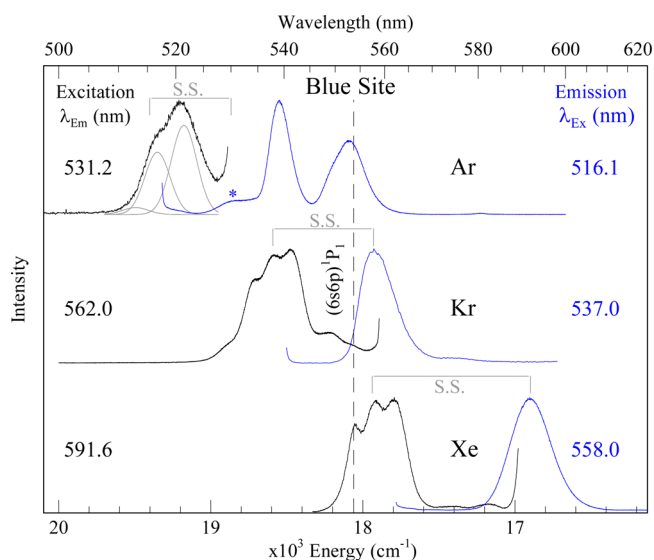


Figure 7. Summary of site-selected excitation (black) and emission spectra (blue traces) recorded for Ba atoms isolated in the blue site in the three rare gas solids studied in the present work. Threefold splitting is very evident in the Xe and Kr excitation spectra and not so in Ar; the presence of three components in the blue site excitation spectrum in this solid was revealed in the Gaussian fit shown by the gray solid traces. The location of the weak blue-site emission in Ba/Ar, due to the low occurrence of this site, is indicated by the asterisk. The horizontal lines labeled “S.S.” show the excitation and emission values used to calculate the blue site Stokes shifts in each host. Each RG matrix has been annealed to 34, 45, and 60 K for Ar, Kr, and Xe, respectively.

respect to minimum of the ground state, to a greater extent when compared to the Ba/Ar and Ba/Kr systems.

III.d. Green Site Luminescence. A summary of the luminescence specific to the green site of isolation is provided in Figure 8. The preference for the occupancy of atomic barium in the green site in the three rare gases matrices is the opposite of that of the blue site. In Ar, the green site is the dominant trapping environment. In Kr, this dominance is matched by both the violet and blue sites, while in Xe, the absorption of the green site is much weaker.

Examination of the green site excitation profiles in Ar and Kr reveals, as highlighted in red in Figure 8, an asymmetric band shape that tails in both cases to higher energy with decreasing intensity. As splitting patterns are not obvious in any of the excitation profiles, the underlying structures of these asymmetric bands were extracted in line shape analyses. Three Gaussian curves (shown in gray in Figure 8) of equal width were required to adequately fit the profiles in Ar and Kr. It is likely that these bands are broadened by the dynamic Jahn–Teller effect and the constituent peaks are unresolved. It is thereby proposed that the site of occupation is isotropic (or at least not far from it), perhaps one of high (cubic) symmetry given the nature of the ground state barium atom. If this is the case, then the green site of isolation must exhibit a larger radius than the blue site in each solid. Inspection of Table 1 reveals that the gas-to-matrix blue shifts for the green site in Ar and Kr are much less than those of the blue site. In addition, the partially resolved band in Xe is more red-shifted than the corresponding blue site in this host. Thus, the confinement of the excited 6p orbitals must be less significant for Ba isolated in the green site. The matrix shifts observed therefore depend

more on the RG polarizability, leading to a red shift of the band center with respect to absorptions of the blue site.

Consistent with the confinement argument just presented, inspection of Table 2 reveals that the calculated Stokes shifts for this site are much less than those recorded for the blue site in each RG host. Moreover, the magnitude of the shift does not increase systematically across the three hosts, with values of 321, 234, and 327 cm⁻¹ measured for Ar, Kr, and Xe, respectively. This is in contrast to the trend observed for the blue site. The green site in Kr shows the smallest Stokes shift and emission bandwidth of the site systems identified in the three hosts. Overall, the emission bandwidths are noticeably narrower than those of the blue site. The combination of small Stokes shifts and narrow emission bandwidths indicates that the equilibrium configuration of the barium atom in the green site is nearly the same in the ground and excited states. Thus, a “vertical” Franck–Condon transition in absorption prepares the excited Ba atom in a configuration which is relatively close to the minimum of the excited state potential energy surface.

III.e. Violet Site Luminescence. A summary of the Ba/RG violet site luminescence is provided in Figure 9, which reveals a splitting pattern in excitation that is in stark contrast to the familiar threefold split JT pattern exhibited on P ← S transitions of most matrix-isolated atoms. The classical JT structure is most clearly presented in the Ba/RG system by the blue site. Contrasting with this, the excitation profiles of the Ba/RG violet sites display an asymmetric 2 + 1 splitting pattern, with the doublet at higher energy and a lower energy singlet. Each component peak of this 2 + 1 structure is clearly identifiable in the Ba/Xe system as shown in the bottom of Figure 9. In contrast, only the higher energy doublets are immediately evident in Ar and Kr. For both Ba/Ar and Ba/Kr the singlet band is strongly overlapped by the green site absorption band. As a result, the location of the singlet component was identified in the line shape fits presented by the gray traces in Figure 9. The singlet band centers are shown by the short red vertical lines.

Examination of Figure 9 and Table 1 indicates that this site of isolation presents two very different interactions for the Ba atom excited into its 6s6p configuration. The splitting between the doublet and singlet features is substantial in each solid. Values of 794, 711, and 633 cm⁻¹ are determined (from the lowest energy singlet to the center of the doublet) in Ar, Kr, and Xe, respectively. This data may be interpreted by considering the orientation of the 6p_x, 6p_y, and 6p_z orbitals of the excited Ba atom in the environment of the violet site. Evidently the interaction in this environment is anisotropic, such that two of the p orbitals (p_x and p_y) are pointing toward one or more RG atoms (“equatorial”, Π configuration) and thereby experience a strong repulsion. It is of note that the interaction of these two orbitals is slightly different, giving rise to a nondegenerate Π doublet in absorption. Contrasting this, the p_z orbital must be directed toward more empty space (“polar”, Σ configuration) and experiences less repulsion with the host atoms. Thus, the higher energy doublet corresponds to the more repulsive Π configuration and the lower energy singlet represents the Σ configuration. Such a Σ–Π splitting was also observed for the excited Ba(¹P) atom on the surface of an Ar cluster.¹³ However, in that case, the Π doublet was lower in energy than the Σ singlet. The 6p_z orbital of Ba was directed toward the Ar cluster and 6p_x and 6p_y orbitals were aligned parallel to the cluster surface, pointing toward empty space. The asymmetric threefold (2 + 1) splitting pattern

Table 2. Photophysical Characteristics of Ba/RG Sites of Isolation As Revealed by $(6s6p) \ ^1P_1 \rightarrow \ ^1S_0 (6s^2)$ Emission Spectroscopy^a

Sr/RG site	λ_{ex} (nm)	λ_{em} (nm)	ν (cm^{-1})	δ (cm^{-1})	SS (cm^{-1})	$\Delta\nu$ (cm^{-1})
argon						
violet	510.3				1047	
	515.4	539.1	18 549	489	853	163
	535.0				142	
blue	516.3	531.2	18 825	765	543	225 ^b
green	528.3	537.4	18 608	548	321	183
red	523.9				902	
	527.7	549.9	18 185	125	765	163
	543.8				204	
krypton						
violet	529.1				921	
	533.4	556.2	17 979	-81	769	129
	552.7				114	
blue	537.9	557.9	17 924	-136	666	261
green	547.8	554.9	18 021	-39	234	135
red ^c	541.4				828	
	544.3	567.0	17 643	-417	729	200 ^b
	561.7				160	
xenon						
violet	549.3				907	
	553.9	578.1	17 298	-762	756	213
	572.0				184	
blue	557.8	591.6	16 903	-1157	1024	311
green	566.4	577.1	17 328	-732	327	202

^a δ represents the gas phase to RG matrix shift calculated from the gas phase value¹⁰ for this transition at 553.7 nm (18 060.261 cm^{-1}). SS represents the Stokes shift. For the blue and green sites, this value was determined from the center of the emission band to the center of the excitation band. For the violet and red sites, three values are quoted which were evaluated from the center of the emission band to each excitation component of the asymmetric threefold split. The last column, $\Delta\nu$, gives the full width at half-maximum values for each emission band in units of wavenumber (cm^{-1}). The wavelength uncertainties in the emission values are estimated to be <0.06 nm. In this spectral range this translates into an uncertainty of ± 2 cm^{-1} . ^bThe spectral bandwidths of the blue site emission in Ar and the R1 site in Kr were estimated with a line shape analysis as significant overlap with emission from other sites occurred in both cases. ^cThermally unstable site of isolation.

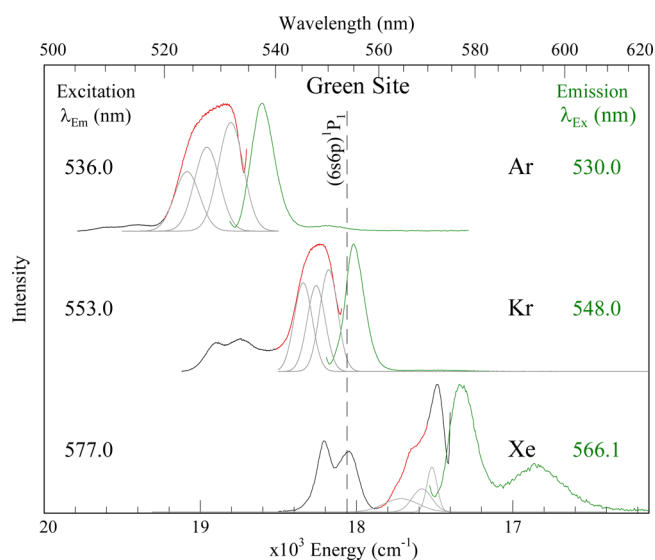


Figure 8. Summary of excitation (red) and emission spectra (green traces) recorded for Ba isolated in the green site of the RG solids. The three Gaussian curves used to fit the green site excitation band in Ar and Kr are shown by the solid gray traces. For clarity, the portion of each excitation band associated with the green site of isolation is highlighted in red. This band is especially weak in Ba/Xe due to the low occurrence of this site in xenon.

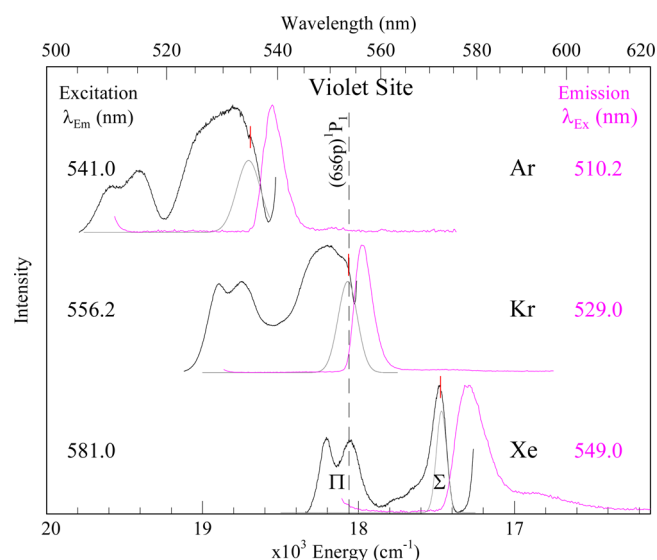


Figure 9. Summary of excitation (black) and emission spectra (purple traces) recorded for atomic barium isolated in the violet site in each RG solid. The Gaussian peaks (shown in gray) indicate the position of the lowest energy component of the asymmetric 2 + 1 split band. Although the position of this band is obvious in Xe, the corresponding components in Ar and Kr are obscured by an intense green site absorption. For clarity the position of this band is shown by vertical red lines for these two solids.

observed for Ba in the RG solids suggests that the violet site of isolation corresponds to a highly axial site. Such a vacancy could be described in a generic sense as a “cylindrical” site. The fcc void that most closely matches this description is a divacancy, but as was found in the recent molecular dynamics (MD) calculations¹ on Ba/Xe, this site is too small even in xenon to accommodate the large $6s^2$ metal atom. The issue of site attributions will be made in section III.h.

In all three hosts, rapid (subnanosecond) nonradiative relaxation occurs to the lowest energy Σ configuration, as only a single emission band is observed regardless of which of the $2 + 1$ components is photoexcited. The details of the Stokes shifts computed from the violet site Π and Σ bands in each host are presented in Table 2. Examination of these values demonstrates that a much larger Stokes shift occurs with excitation of the Π bands—more than 8 times the value of the Σ Stokes shift in the case of Kr. Thus, following excitation of an electron into one of the Π orbitals, a reorientation occurs along the excited state surface to minimize the initial strongly repulsive interactions and thereby lower the energy. From the Σ orientation, the SS values are small, having values of only 142, 114, and 184 cm^{-1} respectively in Ar, Kr, and Xe. The violet site excited state lifetime measurements conducted in Xe show the same nanosecond decay profile, regardless of the p-orbital orientation reached in the excited state. This behavior indicates that orbital reorientation occurs on a subnanosecond time scale.

III.f. Red Site Luminescence. A summary of the red site luminescence features in Ar and Kr is presented in Figure 10, and the photophysical characteristics of this site are collected in Tables 1 and 2. A detailed analysis, which included measuring emissions lifetimes and their temperature dependence, was conducted for Ba/Ar.¹² The same extensive study

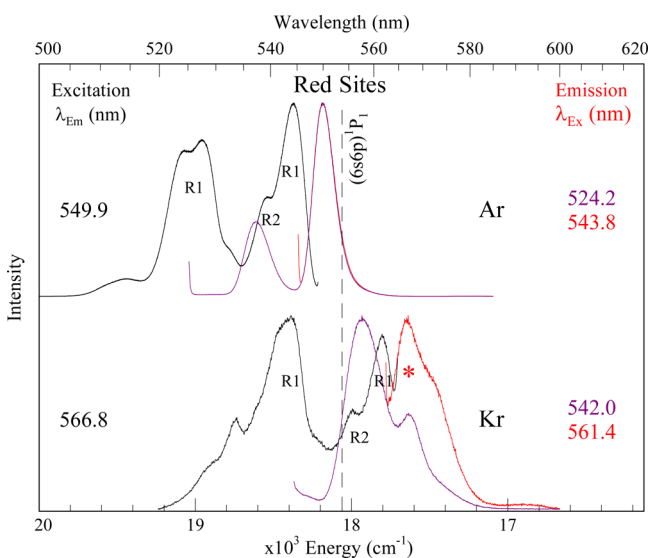


Figure 10. Summary of $(6s6p) \ ^1P_1$ state luminescence of Ba isolated in the red site in solid Ar and Kr samples which were annealed to 28 and 38 K, respectively. The red site in Kr was shown to be thermally unstable with annealing to 45 K. The same site in Ar can be removed by annealing to 34 K—with a substantial loss of the sample. As indicated in the upper portion of the figure, excitation into R1 at two different locations (543.8 and 524 nm) produces the same emission centered at 549.9 nm. In contrast, excitation into R2 at 538.9 nm produces distinct emission centered at 549.0 nm with a larger bandwidth.

could not, however, be achieved for the equivalent bands in Kr owing to the thermal instability of this site in Kr. Moreover, numerous distinct lower energy bands complicate the 2D-EE spectrum (see Figure 4) of a Kr sample annealed to 38 K—indicating that a multitude of partially stable sites exist. For comparative purposes the excitation profile of the brightest emission in this region is presented in the bottom panel of Figure 10.

Even though the red-site emission in Ba/Ar appears simple, its excitation profile is complex. From carefully monitoring the emission and after conducting fits of the excitation profile, it is proposed to consist of two distinct matrix sites, designated R1 and R2, and whose locations are indicated in Figure 10. Line shape fits of the Ba/Ar red site excitation profile required the use of six Gaussian functions, the results of which are shown in Figure 11. Even though only five features are evident in the

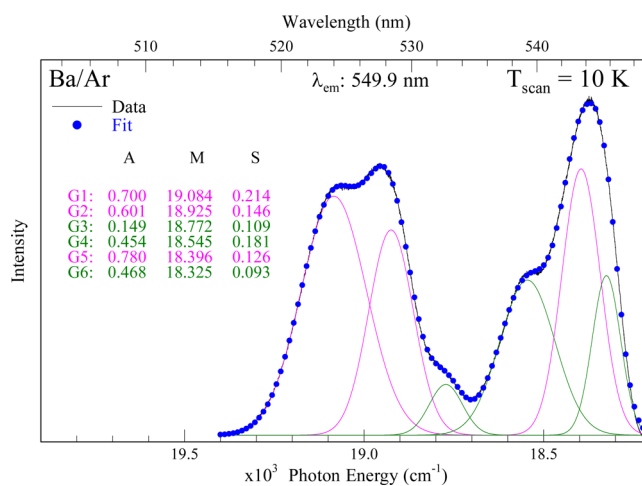


Figure 11. Line shape analysis conducted on the dominant, red site of isolation of atomic barium isolated in solid argon. Even though only five features are evident in the red site excitation profile, acceptable fits required the use of six Gaussian functions (G_n). The parameters of the six functions are listed as A, M, and S—the amplitude, band position, and bandwidth ($\times 10^3 \text{ cm}^{-1}$), respectively. The six curves are color coded revealing the presence of an asymmetric $2 + 1$, R1 (pink) band and a symmetric threefold split, R2 (green) band. The position of the lowest energy component of the R2 band is not obvious in the recorded profile but is essential to obtain an acceptable fit. On the basis of the fit shown, the R1 feature has a symmetric $2 + 1$ splitting while the R2 is a more standard Jahn–Teller structured band.

recorded excitation profile, the best fits consistently indicated the existence of an additional, unresolved band (the lowest energy component) at $18\,325 \text{ cm}^{-1}$. On the basis of the spectral locations and their relative intensities, the six components are grouped as two sets of bands—consisting of the asymmetric $2 + 1$ bands (R1 feature) and a more typical, JT threefold-split (R2) band. The latter extensively overlaps the lowest energy component of the $2 + 1$ band so that a clear distinction of their emission cannot be provided. However, as indicated in the upper portion of Figure 11, excitation into R1 at two different locations (543.8 and 524 nm) produces the same emission centered at 549.9 nm. In contrast, excitation into R2 at 538.9 nm produces distinct emission centered at 549.0 nm with a larger bandwidth.

Analysis of the luminescence in Ar suggests that the primary contributor is a trapping site (R1) which gives rise to a “doublet plus singlet” splitting pattern, similar to the violet site.

A comparison with the thermally unstable red site profile of Ba/Kr in Figure 10 (lower panel) reveals the same general features. The equivalent of the Ar higher energy doublet centered at ~ 543 nm is not resolved in Kr but does exhibit a dominant peak at 544.3 nm and a shoulder at 541.4 nm. The band at 555.8 in Kr is tentatively assigned as a secondary, minor red site (R2) based on the distinct emission produced at this excitation wavelength, as can be seen in the 2D-EE spectrum presented in Figure 4 (left panel).

On the basis of this interpretation, the primary R1 site of isolation gives rise to two distinct interactions for the Ba atom excited into the $(6s6p) \ ^1P_1$ state. A large splitting between the doublet and singlet feature is observed in Ar and Kr. Values of 699 and 619 cm^{-1} are determined (from the lowest energy singlet to the center of the doublet) in Ar and Kr respectively—noticeably smaller than those obtained for the violet site. The splitting may be explained by considering the orientation of the $6p_x$, $6p_y$, and $6p_z$ orbitals of the excited Ba atom with respect to the anisotropic configuration of the surrounding RG atoms. Similar to the violet site, the higher energy doublet corresponds to the more repulsive Π configuration (p_x and p_y) and the lower energy singlet represents the Σ configuration (p_z). It is interesting to note that the extent of splitting on the Π doublet is less than that of the violet site. Combining this characteristic with a smaller overall singlet to doublet energy gap suggests a weaker interaction of the excited Ba atom at this site in comparison to the violet site. A more voluminous cylindrical site could therefore be proposed for R1. As the features of R1 are connected to a single emission band in Ar and Kr, a relaxation to the lowest energy excited $(6s6p) \ ^1P_1$ state surface must occur in both hosts.

III.g. Temporal Analysis. Time-resolved emission spectra (TRES) were recorded, as described in the Methods section and elsewhere,^{1,12} with nanosecond pulsed laser excitation and time-gated iCCD detection. The decay profiles extracted were, in most cases, well-fitted with single exponential functions—typical data and fits were shown in the recent Ba/Xe study.¹ This allowed the excited state lifetimes to be determined, and the observed values (τ_{obs}) for each site of isolation in the three Ba/RG systems are collected in Table 3. With the exception of the blue site in Ba/Xe, all measured decay times were found to be invariant over a temperature range of at least 3 K, indicating that the radiative lifetimes had been identified. The lifetime values are plotted in Figure 12 along with a λ^3 extrapolation of the gas phase¹⁰ value of 8.4 ns. As shown, the directly measured (τ_{obs}) matrix values are all less than the gas phase extrapolation shown by the solid curve. However, in a comparison with the gas phase values, a correction must be made for the effective field of the solid.⁸ As described in a previous study on the Zn/RG systems,¹⁴ this leads to an enhancement of electronic transitions and a shortening of the observed matrix lifetimes. The correction^{15,16} for the effective field of the solid host is made using the following equation:

$$\tau_{\text{cor}} = \tau_{\text{obs}} n [s(n^2 - 1) + 1]^2 \quad (1)$$

Equation 1 is a function of the index of refraction,¹⁷ n , of the rare gas host and a “shape” parameter, s . The latter is a depolarization factor related to the symmetry of the site accommodating the guest atom. For an isotropic, spherical cavity $s = 1/3$. However, this value may not be appropriate for all trapping environments of Ba in the solid RGs, i.e., the violet and red sites. The asymmetric threefold splits observed for

Table 3. Summary of Observed and Corrected Excited State Lifetimes for Each Site of Isolation in All Three RG Systems^a

RG	site	τ_{obs} (ns)	τ_{rad} (ns)	τ_{cor} (ns)			
				$s = 0$	$s = 1/3$	$s = 1/2$	$s = 1$
Ar	violet	4.5	✓	5.9	9.2	11.2	18.0
	blue	3.9 ^b	✓ ^b	5.1	8.0	9.7	15.6
	green	4.3	✓	5.7	8.8	10.7	17.2
	red	4.9	✓	6.5	10.1	12.2	19.6
Kr	violet	4.2	✓	6.0	10.9	13.8	24.9
	blue	4.4	✓	6.3	11.4	14.5	26.1
	green	4.5	✓	6.4	11.6	14.8	26.7
Xe	violet	4.3	✓	6.4	12.7	16.6	31.6
	blue	3.3	×	4.9	9.7	12.7	24.2
	green	4.0	✓	6.0	11.8	15.4	29.4

^aThe effective field correction was made with eq 1. With the use of single exponential fits, the uncertainty in the extracted decay times is estimated to be ± 0.2 ns. It arises from a convolution of the 5 ns pulsed laser excitation and the 2 ns iCCD detection. ^bThe decay time obtained for the blue site in Ar is only tentatively assigned as the radiative lifetime.

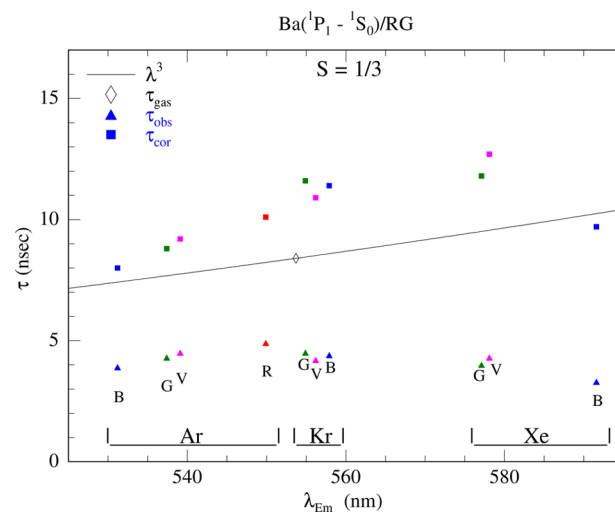


Figure 12. Comparison of emission decay times recorded at 10 K for the thermally stable Ba: $(6s6p) \ ^1P_1 \rightarrow \ ^1S_0 (6s^2)$ emission features in Ar, Kr, and Xe and a λ^3 extrapolation of the gas phase lifetime¹⁰ for this transition of atomic Ba. With the sole exception of the blue site in Ba/Xe, all the measured decay times were found to be independent of temperature and accordingly are identified as radiative lifetimes. The uncorrected lifetime values are given by filled triangles, whereas the values corrected for the effective field of the host are denoted by the filled squares.

these sites implies that the RG site accommodating the large Ba atom is anisotropic. As a test, a number of values for s were used to compute the effective field correction based on a variety of cavity shapes, as described by Shibuya.¹⁵ For instance, transitions whose moments are considered to be in the longitudinal or transverse direction of a long cylindrical cavity yield values of $s = 0$ and $1/2$, respectively.

Inspection of the corrected lifetime values provided in Table 3 shows that the values closest to the gas phase λ^3 extrapolation (Figure 12) are those computed with the parameter $s = 1/3$. However, even these show a deviation and are all, with the exception of the blue site in Ba/Xe, systematically longer than the gas phase extrapolation. It

should be noted that the radiative lifetime of this site in xenon was not identified even at 9.8 K, the lowest temperature attainable with the cryostat used, so the recorded value does not provide a valid comparison with the other sites. This model, using just the index of refraction of the solid, is probably too simplistic to take into account the complicated interactions that occur for an excited Ba atom at a given RG site and the effects such an interaction has on the oscillator strength of the transition involved. Nonetheless, the measured nanosecond radiative decay times allows confident assignment of the observed Ba/RG emission bands to resonance fluorescence of the $(6s6p) \ ^1P_1 \rightarrow \ ^1S_0 (6s^2)$ transition of atomic Ba.

III.h. Site Attributions. The excitation profiles recorded for three Ba/RG matrix systems display the richest variety of structures for any matrix-isolated metal atom studied to date. The structures range from the classic Jahn–Teller threefold split bands—best exhibited by the blue site—to asymmetric threefold $(2 + 1)$ patterns on the violet site, to the apparently structureless bands presented by the green site. The violet, $2 + 1$ bands are observed for all three hosts examined in the present study and, to the best of our knowledge, are unique to the Ba/RG matrices. However, from the linear dependence observed in the matrix shift vs matrix polarizability plots, each of these colored sites corresponds to a particular site in the three hosts.

Moreover, on the basis of the results obtained from recent molecular dynamics calculations done on the Ba/Xe system, the nature of the sites producing the three main spectral features—violet, blue, and green—can be attributed to particular site types. In the Ba/Xe study, strong evidence emerged from comparisons of the recorded excitation spectra with spectral simulations that the blue site was due to a symmetrical tetravacancy (TV) site and the green site was due to a hexavacancy (HV) site in the crystalline fcc lattice. These Ba/Xe predictions are consistent with the pronounced threefold splitting observed on the blue sites in Kr and Xe and also with the three component fits done on the green site for all three solids. Thus, in the cuboctahedral fcc lattice, the TV and HV multivacancies result in high symmetry O_h and T_d sites producing the characteristic Jahn–Teller splitting on the isolated guest metal atom. More uniquely, the molecular dynamics calculations indicated that the trapping site responsible for the $2 + 1$ pattern of the Ba/Xe violet site corresponds to a Ba atom isolated at a grain boundary (GB). In solid xenon this site consists of a five-atom vacancy with D_{3h} symmetry, located at the boundary of two fcc slabs. From the similar patterns observed for the violet site in Kr and Ar and the linear matrix shift vs polarizability plots, it is highly likely that this site is also present in both of these solids.

Attributions of the thermally stable red sites that exist in Ba/Ar and the partially stable red sites in Ba/Kr are less clear. While a large variety of multivacancy sites can exist at grain boundaries, Monte Carlo simulations conducted on Na atoms in solid Ar¹⁸ showed that matrix sites possessing a strong permanent axial asymmetry, such as a cylindrical vacancy, yielded a widely split “doublet plus singlet” absorption line shape. The R1 trapping sites in Ar and Kr which give rise to the $2 + 1$ splitting of the excited P state of Ba must be characterized by a lower symmetry since the doublet component exhibits a small splitting. Recent molecular dynamics simulations, conducted by Gervais et al.¹⁹ of the alkali metal atoms embedded in Ar, may offer insight into such

vacancies. The study revealed that the “red” absorption bands of large atoms, such as potassium, arise from the occupation of a number of distinct grain boundary (GB) sites in Ar. Indeed these sites were shown to be more favorable than larger vacancies in the bulk solid. The authors report that more compact GB cavities with reduced dimensionality systematically exhibit one or more peaks shifted toward higher energy. It is plausible that such a vacancy could reproduce the excitation/absorption patterns intrinsic to the occupation of a Ba atom in the R1 sites. On the other hand, larger GB cavities can produce broad unstructured bands and triplet structures (not arising from a JT effect)—which could be correlated to the R2 matrix site.

However, the GB sites are known¹⁹ to be quite sensitive to matrix annealing. This behavior closely matches the relatively low annealing threshold observed for the red (R1 and R2) site features in Ar and Kr but is in conflict with the high stability of the violet sites in all three hosts. From the existing theoretical work done on Ba/Xe, it is proposed that the violet sites of isolation occur at a GB vacancy that involves the minimum loss of structure within the bulk RG solid. For Ba/Xe, a compact five-atom vacancy which maintains threefold symmetry of the fcc lattice was identified as the origin of the violet band. Based on the observed matrix shifts, the red excitation features (R1) in Ar and Kr should correspond to the isolation of Ba in larger GB cavities. The exact nature of the red sites requires input from MD calculations and spectral simulations. It is hoped future work should be revealing of this dominant site in Ba/Ar.

IV. CONCLUSION

The 2D-EE spectra recorded in this study have allowed a comprehensive yet detailed examination to be conducted of the sites occupied by Ba atoms in the three solid rare gases Ar, Kr, and Xe. The excitation scans extracted from the 2D-EE spectra reveal that the site occupancy of barium in the RG solids is much more complex than the absorption spectroscopy would suggest. Three thermally stable sites—labeled blue, green, and violet—have been identified in all three solids. The amounts of these sites vary from solid to solid—with the blue site dominating in Xe while it is the green in Ar. In Ba/Kr the three sites are present in roughly similar quantities.

The temperature dependence of the excited state lifetimes was found to depend primarily on the RG host, but also on the site of isolation. For Ba/Xe, a shortening of the site-specific emission lifetimes at elevated temperatures was observed¹ and tentatively attributed a $^1P_1 \rightarrow \ ^3P_2$ intersystem crossing. This effect was most pronounced for the blue site of isolation, where the radiative lifetime was not identified, even at 10 K. In contrast, changes in the decay profiles were not observed in Ar and Kr, over the temperature ranges assessed. The efficiency of intersystem crossing in solid Xe is proposed to arise from the external heavy atom effect.⁸

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Notes

The authors declare no competing financial interest.

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